NAVY UNDERWATER SOUND LAB NEW LCHDON CONN
SIGNAL-TO-NOISE GAIN OF AN IDEALIZED ARRAY PATTERN AS A FUNCTIO--ETC(U)
APR 62 M C KARAMARGIN, B F CRON
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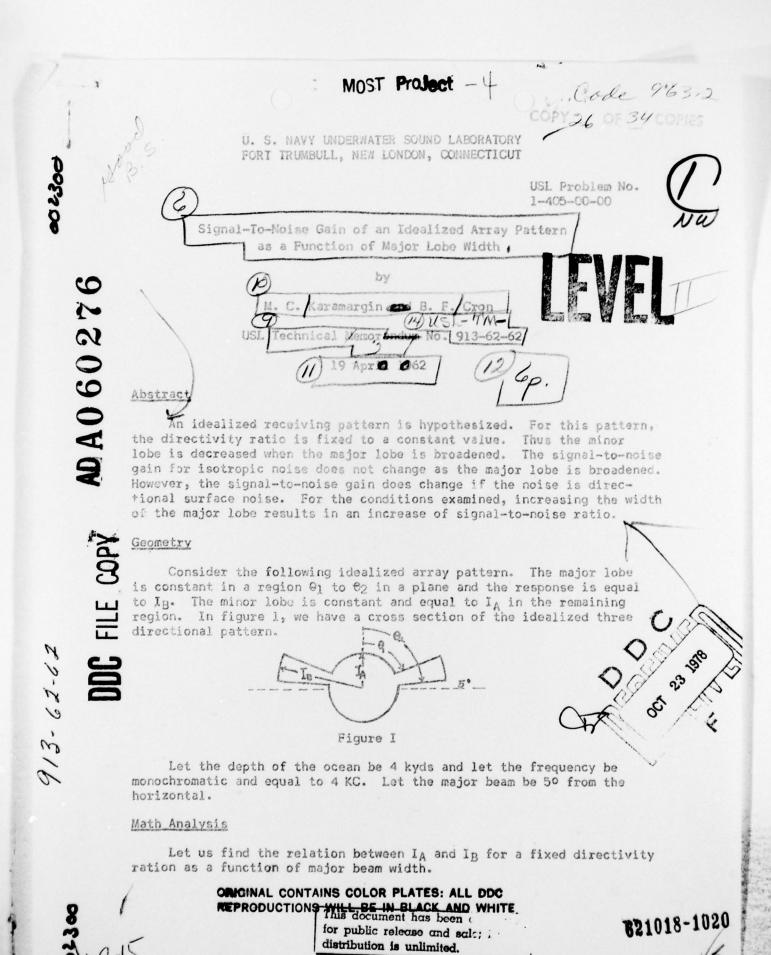












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D.R. = $I_{\rm B/I_{AV}}$ where $I_{\rm B}$ is the maximum intensity received and $I_{\rm AV}$ is the intensity averaged over all directions. Let us draw an arbitrary sphere of radius V about the center of the array and consider spherical coordinates. Then

Let us set D.R. = K

Then

$$I_{A} = \frac{I_{B}\left[1 - \frac{K}{2}(\cos\theta_{1} - \cos\theta_{2})\right]}{K\left[1 + \frac{1}{2}\left(\cos\theta_{2} - \cos\theta_{1}\right)\right]}$$

Thus for major beam widths of 5° , 10° and 15° (Θ_2 - Θ_i), and for a given K, we can find the relationship between I_A and I_B . For all these cases, the signal-to-noise gain will be the same if the noise is isotropic. We will now investigate the signal-to-noise gain for the case of directional sufrace noise.

Surface Noise

Consider surface noise which is directional in the direction and radiates independently in the azimuthal direction. We will consider attenuation and spherical spreading. Let the directionality of the intensity of the radiated source be $\cos^m\Theta$. Then the intensity received at the array due to one noise source is

where $\cos^{\infty} \Theta$ is the directional property of noise, $\exp - \infty r$ is the attenuation factor, r^2 is the spherical spreading factor.

oc = 0.17 in 10

Let us now assume that there is one noise source per unit area. Then in a ring of noise sources the energy received on the major beam is

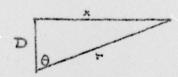
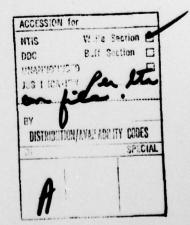


Figure 2



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By changing variables

Is received =
$$2\pi(\alpha D)^{3n} \int \frac{\exp(-y)}{y^{m+1}} dy$$
 In received = $2\pi \operatorname{In}(\alpha D)^{n} \int \frac{\exp(-y)}{y^{m+1}} dy$

where

The region from $\Theta = \Theta_2$ to $\Theta = \frac{\pi}{2}$ has been disregarded since it will be the same form in all three cases of comparison.

Using a directivity ratio of K=10 and varying the beam width, it is found that as the major lobe increases, the total received noise decreases. The same is true for a K=5. The cos Θ and $\cos^2 \Theta$ directionality of the noise is compared.

It is interesting to compute the received intensity as a function of direction. For a $\cos^{\frac{1}{2}}\Theta$ directionality of the noise, the received intensity in the angle $40~\text{Å}\phi$ is

For $\cos \theta$ and $\cos^2 \theta$ noise sources, the received intensity is given in graph 2.

Conclusions

For the idealized conditions shown, increasing the major beam width results in a lower noise input and in a larger signal-to-noise ratio. Since this increase is on the order of 18 db, it is worthwhile computing the gains for more realistic situations. From graph 2 we see the received intensity as a function of direction. For this type of situation, the minor lobes should be placed at the maximum intensity and the major lobe at the small received intensity.

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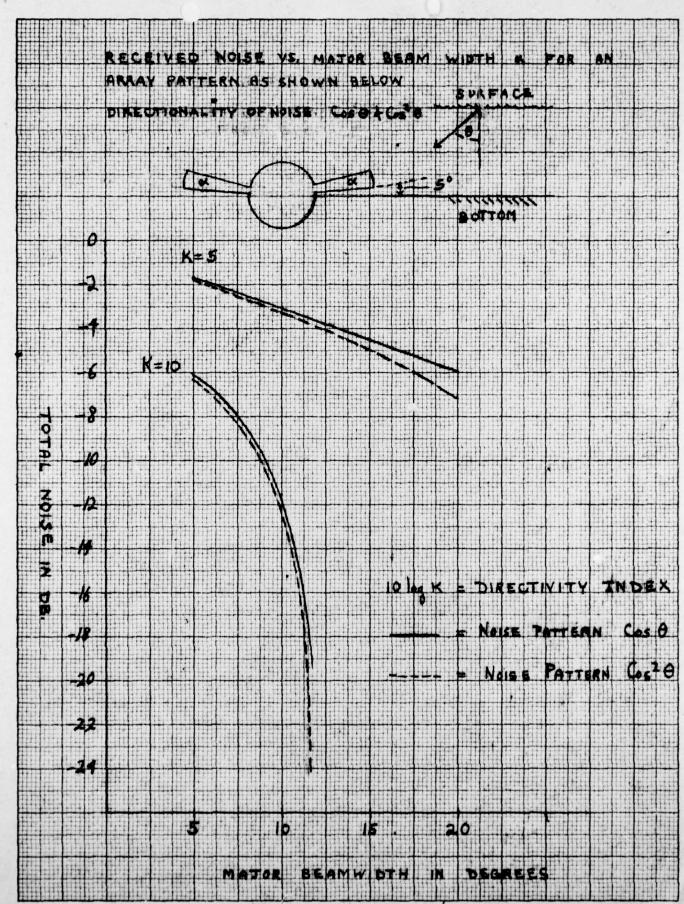
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REFERENCES

I. R. J. Urick - "Some Directional Properties of Deep-Water Ambient Noise" NRL Report 3796 dtd 16 Jan. 1951

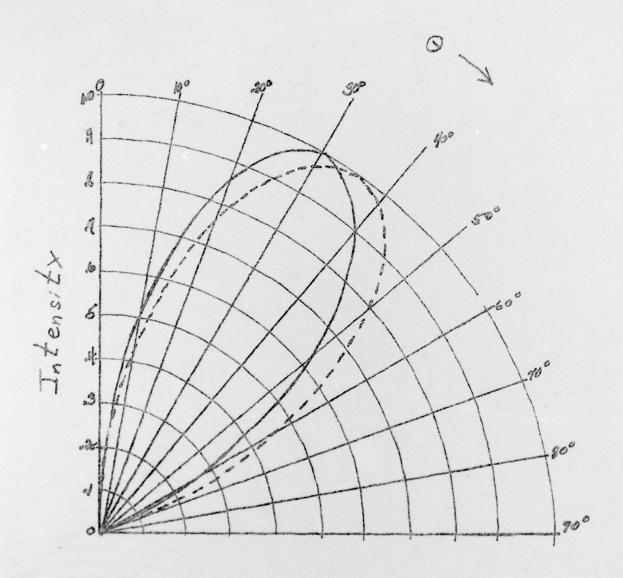
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GRAPH I

--- Coo 9 --- Coo 3



Received Intensity in an Angle dodo As a Function of Direction Graph 3.